

# Energy Levels of Light Nuclei

## $A = 10$

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**Abstract:** An evaluation of  $A = 5\text{--}10$  was published in *Nuclear Physics* 78 (1966), p. 1. This version of  $A = 10$  differs from the published version in that we have corrected some errors discovered after the article went to press. Figures and introductory tables have been omitted from this manuscript. Reference key numbers have been changed to the TUNL/NNDC format.

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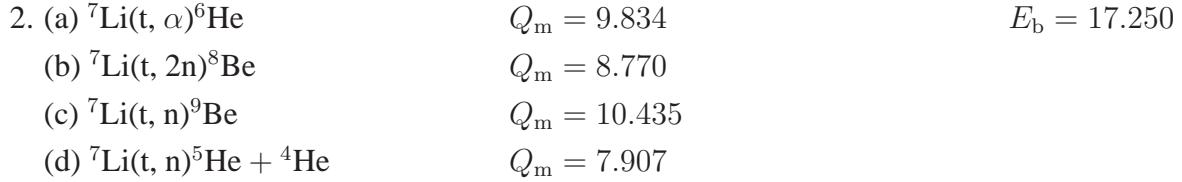
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<sup>10</sup>**Be**  
(Figs. 19 and 22)

GENERAL: See ([1956KU1A](#), [1957FR1B](#), [1959BA1F](#), [1960KU1B](#), [1960TA1C](#), [1961BA1E](#), [1961TR1B](#), [1963BU1C](#), [1963VL1A](#), [1963WA1M](#), [1964FR1D](#), [1964GR1J](#), [1964VO1B](#), [1964WA1F](#), [1964WA1K](#)).



The weighted mean end-point energy is  $0.556 \pm 0.003$  MeV ([1951LI26](#)). The mean half life is  $(2.7 \pm 0.4) \times 10^6$  y ([1949HU19](#)):  $\log ft = 13.65$  ([1951FE1A](#)). The spectrum is of the D<sub>2</sub> type ([1950WU1A](#)).



The neutron yield exhibits broad resonances at  $E_t = 0.84$  and  $1.70$  MeV ([1951CR01](#)),  $0.71 \pm 0.02$  MeV ([1962SE1A](#)),  $0.765$  and  $1.735$  MeV ([1961VA43](#)), and a weak structure at  $E_t = 0.24$  MeV ([1960SE12](#)). The width of the  $0.77$  MeV resonance is  $160 \pm 50$  keV ([1962SE1A](#)).

Excitation functions for  $\alpha_0$ ,  $\alpha_1$  and  $\alpha_2$  (corresponding to  $^6\text{He}^*(0, 1.8, 3.4)$ ) see, however, ([1965AJ01](#)) for evidence relevant to the existence of a state at  $3.4$  MeV in  $^6\text{He}$ ) show a weak maximum at  $1.10$  MeV, a weak minimum at  $1.30$  MeV and a broad maximum at  $1.80$  MeV ( $\Gamma \approx 0.5$  MeV). Angular distributions are complex and suggest direct interaction: the maximum at  $E_t = 1.80$  MeV may result from interference effects ([1961HO23](#)). Angular distributions of  $\alpha_0$  and  $\alpha_1$  at  $E_t = 0.24$  MeV suggest  $J = 2^+$  for the  $17.8$  MeV state ([1954AL38](#)). See also ([1963JA1E](#)) and ([1959AJ76](#)).



Angular distributions of  $p_0$  at  $E_\alpha = 13.6$  and  $14.7$  MeV ([1962KO13](#)) and  $30$  MeV ([1960KL03](#)) show an oscillatory character with strong peaking in the back hemisphere. The  $p_1$  distributions are smoother, but they also show backward peaks (see also ([1956WA29](#))). An explanation in terms of heavy-particle stripping is discussed by ([1962HO1C](#)). See also ([1962MA59](#)) and ([1960MA15](#)).



Table 10.1: Energy levels of  $^{10}\text{Be}$ 

$E_x$ (MeV $\pm$ keV)	$J^\pi$	$\Gamma$ (keV)	Decay	Reactions
g.s.	$0^+$	$\tau_{1/2} = (2.7 \pm 0.4) \times 10^6 \text{ y}$	$\beta^-$	1, 3, 4, 5, 11, 14, 17, 19, 21
$3.366 \pm 3$	$2^+$	$\tau_m = 0.15 \pm 0.03 \text{ psec}$	$\gamma$	3, 4, 5, 11, 17
$5.959 \pm 5^{\text{a}}$	$1^-, 2^-$	$\tau_m < 0.08 \text{ psec}$	$\gamma$	4, 5, 6, 11
$6.178 \pm 9$	$0^+$	$\tau_m > 0.5 \text{ psec}$	$\pi, \gamma$	4, 11
$6.262 \pm 9$	$2^-$		$\gamma$	4, 11
$7.377 \pm 10$	$3^-$	16	n	4, 6, 11
$7.548 \pm 10$	$2^+$	6	n	4, 6, 11
9.27	$(4^-)$	100	n	4, 6, 10
(9.4)	$(2^+)$	$\approx 400$	n	4, 6, 10, 17
10.7	$\geq 1$		$n, \alpha$	4, 6, 17
(17.42)			n, t	2
17.79	$(2^+)$	$110 \pm 35$	$n, t, \alpha$	2
18.47		$\approx 500$	$n, t, \alpha$	2
24				17

<sup>a</sup> See [Note added in proof](#) section in the Introduction here.

See (1964CA16).

5.  ${}^9\text{Be}(n, \gamma){}^{10}\text{Be}$   $Q_m = 6.815$

The thermal capture cross section is  $9.5 \pm 1.0 \text{ mb}$  (1964ST25). Reported  $\gamma$ -transitions are listed in Table 10.2 (1960BA01, 1961JA19, 1963DR02).

6.  ${}^9\text{Be}(n, n){}^9\text{Be}$   $E_b = 6.815$

The total cross section data is summarized in (1964ST25). Angular distributions are summarized in (1963GO1M). The coherent scattering length (thermal, bound) is 7.7 fm (1961WI1A). The spin dependent thermal cross section is  $< 30 \text{ mb}$  (1952PA1A).

In the region  $E_n = 0$  to 16 MeV, four resonances are reported, at 0.62, 0.81, 2.73 and 4.3 MeV: see Table 10.3.

Table 10.2: Neutron capture gamma rays in  $^{10}\text{Be}$ 

$E_\gamma$ (MeV $\pm$ keV)	Transition	Intensities <sup>a</sup>		
		A	B	C
6.807 $\pm$ 7	capt. $\rightarrow$ g.s.	62	70	65
5.956 $\pm$ 6	5.96 $\rightarrow$ g.s.	1.4	2	$\lesssim 2$
3.441 $\pm$ 3	capt. $\rightarrow$ 3.37	11	15	11
3.365 $\pm$ 3	3.37 $\rightarrow$ g.s.	37	28	28
2.590 $\pm$ 3	5.96 $\rightarrow$ 3.37	28	17	21
0.855 $\pm$ 3	capt. $\rightarrow$ 5.96	29.4	16	24

A: (([1960BA01](#)): see ([1963DR02](#))).

B: ([1961JA19](#)).

C: ([1963DR02](#)).

<sup>a</sup> Gamma rays per 100 captures.

 Table 10.3: Resonances in  $^9\text{Be}(n, n)^9\text{Be}$ 

$E_{\text{res}}$ <sup>a</sup> (MeV $\pm$ keV)	$^{10}\text{Be}^*$ (MeV)	$\Gamma_{\text{c.m.}}$ (keV)	$J^\pi$	$l$	$R$ (fm)	$\gamma_n^2$ (keV)		$\theta^2$ %	References
						$J_c = 1$	$J_c = 2$		
0.625 $\pm$ 10	7.377	16	$3^-$	2	5.6	82	82	7.5	( <a href="#">1951BO45</a> , <a href="#">1964LA04</a> )
0.815 $\pm$ 10	7.548	6	$2^+$	1	5.6	0.8	5.3	0.28	( <a href="#">1951BO45</a> , <a href="#">1955WI25</a> , <a href="#">1964LA04</a> )
2.73	9.27	$\approx 100$	$(4^-)$						( <a href="#">1951BO45</a> , <a href="#">1959FO1A</a> )
(2.85)	9.4	$\approx 400$	$(2^+)$	(1)					( <a href="#">1951BO45</a> , <a href="#">1959FO1A</a> )
4.3	10.7		$\geq 1$						( <a href="#">1961FO07</a> )

<sup>a</sup> ([1962PE10](#)) report an additional anomaly in the cross section at  $E_n = 207$  keV.

Polarization and differential cross sections are reported for  $E_n = 0.2$  to  $2$  MeV by (1961LA1A, 1962EL01, 1964LA04). Analysis of these data indicates that the  $E_n = 0.62$  MeV ( $E_x = 7.37$  MeV) resonance has  $J^\pi = 3^-$ , formed by  $l = 2$ , with equal participation of channels  $J_c = 1$  and  $2$ . Interference with hard-sphere ( $R = 5.6$  fm) s-wave background and with a broad  $l = 1$ ,  $J = 3^+$  state is observed. Below  $E_n = 0.5$  MeV, the scattering cross section reflects the effect of bound  $1^-$  and  $2^-$  states with  $E_\lambda \approx -0.2$  MeV,  $\gamma^2 \approx 0.05$  MeV, presumably to be identified with  $^{10}\text{Be}^*(5.96, 6.26)$ . The resonant behavior near  $0.815$  MeV is consistent with  $J^\pi = 2^+$ , formed by  $l = 1$  (1964LA04). See also (1964GO1L). The finding that the level at  $E_x = 7.37$  MeV has  $J = 3^-$  is in good accord with comparisons of reduced widths with the presumed  $T = 1$ ,  $T_z = 0$  analogue in  $^{10}\text{B}$  at  $E_x = 8.89$  MeV; for  $R = 5.8$  fm,

$$[\gamma_p^2 + \gamma_n^2](^{10}\text{B})/\gamma_n^2(^{10}\text{Be}) = 1.07 \text{ (1962AL1A).}$$

The structure at  $E_n = 2.73$  MeV is ascribed to two levels: a broad state at about  $2.85$  MeV with  $J = 2^+$ , and a narrow one,  $\Gamma \approx 100$  keV, at  $E_n = 2.73$  MeV with a tentative assignment of  $J = 4^-$  (1951BO45, 1959FO1A). A weak dip near  $4.3$  MeV is ascribed to a level with  $J \geq 1$  (1961FO07).

Elastic scattering studies for  $E_n > 14$  MeV are reported by (1958NA09, 1960MC04, 1960PE25, 1960SA25). See also (1959AJ76). Optical model analyses have been carried out by (1960HO14, 1963LU10).

See also (1960WA07, 1962OT01, 1962PE10, 1963NE1H, 1964CR1B).

7. (a) $^9\text{Be}(n, n')^9\text{Be}^*$	$E_b = 6.815$
(b) $^9\text{Be}(n, 2n)^8\text{Be}$	$Q_m = -1.665$

Data on non-elastic cross sections are summarized by (1964ST25). The processes involved include  $(n, 2n)$  and  $(n, n')^9\text{Be}^* \rightarrow n + ^8\text{Be}$ ; the  $(n, \alpha)$  process is relatively weak. The total non-elastic cross section rises rapidly from threshold at  $E_n \approx 2.5$  MeV to  $\approx 600$  mb at  $E_n = 6$  MeV, and then falls slowly to  $500$  mb at  $14$  MeV. In the range  $3.5$  to  $6.0$  MeV,  $^9\text{Be}(n, n')^9\text{Be}^*(2.43)$  accounts for about half of the non-elastic cross section (1959MA34: see (1964ST25)): see also  $^9\text{Be}$ . At  $E_n = 14$  MeV,  $^8\text{Be}^*(2.9)$  is frequently involved (1959CH1E, 1961MY01).

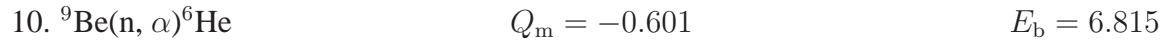
See also  $^9\text{Be}$  and (1959MA1C, 1961CA23, 1961CO1E, 1963DI1F, 1963GO1M, 1963KU1F, 1963MC1C, 1963ZU1B).

8. $^9\text{Be}(n, p)^9\text{Li}$	$Q_m = -12.832$	$E_b = 6.815$
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At  $E_n \approx 15.5$  MeV,  $\sigma(n, p) = 0.7$  mb (1959AL83, 1963AL18, 1964ST25). See also (1960BU1C).

9. $^9\text{Be}(n, t)^7\text{Li}$	$Q_m = -10.435$	$E_b = 6.815$
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At  $E_n = 14$  MeV,  $\sigma(n, t) = 18 \pm 1.5$  mb ([1958WY67](#)). See also ([1958VA33](#)). Measurement of  $\gamma$ -rays from  ${}^9\text{Be}(n, t){}^7\text{Li}^*(0.48)$  gives  $\sigma(n, t_1) = 20, 10,$  and  $30$  mb at  $E_n = 13.6, 14.1$  and  $14.7$  MeV ([1960BE18](#)).



The cross section for production of  ${}^6\text{He}$  has been measured for  $E_n = 0.7$  to  $8.6$  MeV ([1957ST95](#), [1958VA33](#), [1961BA53](#)); see also ([1964ST25](#)). ([1957ST95](#)) find only a smooth rise to a broad maximum of  $104 \pm 7$  mb at  $3.0$  MeV, followed by a gradual decrease to  $70$  mb at  $4.4$  MeV. No indication of resonance is found at  $E_n = 2.7$  MeV. From  $E_n = 3.9$  to  $8.6$  MeV, the cross section decreases smoothly from  $100$  mb to  $32$  mb ([1961BA53](#)). The cross section at  $E_n = 14$  MeV is  $10 \pm 1$  mb ([1953BA04](#)). See also ([1960BU1C](#), [1963CH1C](#), [1964GA11](#)).



Parameters of levels observed in this reaction are listed in Tables [10.4](#) and [10.5](#). Angular distributions of proton groups have been studied at many energies: see ([1958ZE01](#), [1959AJ76](#), [1960MA32](#), [1961IS01](#), [1961RE04](#), [1962SL04](#), [1964SC12](#)). Except at the lowest energies, the stripping process appears to dominate: see ([1959BO1C](#), [1960BE1B](#), [1960LU04](#), [1960LU1B](#), [1960NA1A](#), [1963SM05](#), [1963TA1A](#), [1964BA1V](#)). See also ([1964ZA1B](#)).

The mean life of the  $3.37$  MeV level is  $0.15 \pm 0.03$  psec ([1965WA1P](#)), some ten times shorter than the single-particle value. On the I.P.M. an effective charge  $[1 + (0.5 \pm 0.2)]e$  is required; strong collective effects are indicated ([1963WA03](#), [1963WA1M](#)): see also ([1959BO49](#), [1959KO1B](#)). For the  $6.18$  MeV level,  $\tau_m > 0.5$  psec, consistent with an E0 transition ([1963WA17](#)). Internal pair correlations establish the multipolarities of the transitions:  $3.37$  (E2),  $5.96$  (E1), and  $6.18$  (E0) ([1964WA05](#)); however, ([1965FO1G](#)) report that  $J^\pi(5.96) = 2^-$  from  $\gamma$ - $\gamma$  correlation studies <sup>†</sup>. The mean life of  ${}^{10}\text{Be}^*(5.96)$  is  $< 0.08$  psec ([1965WA1P](#)). The absence of a ground-state transition from  ${}^{10}\text{Be}^*(6.26)$  supports the assignment  $J^\pi = 2^-$  for this level ([1958ME81](#), [1963WA17](#)): see also ([1959CH28](#), [1959GO78](#)).

The (p- $\gamma$ ) correlation through  ${}^{10}\text{Be}^*(3.37)$  has the form  $1 + AP_2(\cos \theta_\gamma)$ , where  $\theta_\gamma$  is measured from the recoil direction. Reported values of  $A$  lie in the range  $-0.2$  to  $-0.4$  ([1957CO54](#), [1959TA01](#), [1959ZA01](#), [1960GO18](#), [1961RE04](#), [1962KO14](#), [1963AC01](#), [1963NE09](#), [1964ZA03](#)). These values are consistent with the assumption of a  ${}^3\text{P}_2$  state in  $LS$  coupling, but not with  ${}^1\text{D}_2$ . Pure  $jj$  coupling leads to  $A = 0$ , as does the collective model. Approximate agreement can be obtained in intermediate coupling with  $a/k = 4.5$  ([1962PI1A](#)): see also ([1960KU1B](#), [1961RO1K](#)). Polarization of the protons has been studied by ([1959HI1E](#), [1960GR11](#), [1960HI09](#), [1961VA03](#), [1962AL10](#), [1962PA12](#), [1963NE09](#), [1963NE16](#), [1964RE04](#), [1965HE1B](#)). See also ([1961TE02](#)).

<sup>†</sup> See *Note added in proof* section in the Introduction here.

Table 10.4: Levels of  $^{10}\text{Be}$  from  $^9\text{Be}(\text{d}, \text{p})^{10}\text{Be}$ <sup>d</sup>

$E_x$ <sup>a</sup> (MeV $\pm$ keV)	$\Gamma$ (keV)	$J^\pi$	$l_n$	$\theta_n^2$ (%)			
				A	B	C	D
0		$0^+$	1	5 – 9	9.3	7.5	9
$3.368 \pm 9$		$2^+$	1	1	1.1	1.1	2.5
$5.959 \pm 9$		$1^-, 2^-$ <sup>b</sup>	0	15	17	12	
$6.178 \pm 9$		$0^+$ <sup>b</sup>		(0.5)			
$6.262 \pm 9$		$2^-$	0	8	18	3.2	
7.37	$\approx 25$	$3^-$	$2^c$	1.3	4.1	2	
7.54	< 10	$2^{(+)}$		0.34			

A: (1958ME81).

B: (1960MA32).

C: (1964SC12).

D: (1961RE04): all PWBA.

<sup>a</sup> (1954JU23).

<sup>b</sup> (1963WA17, 1964WA05, 1965FO1G).

<sup>c</sup> (1964SC12).

<sup>d</sup> See *Note added in proof* section in the Introduction here.

12.  $^9\text{Be}(\text{t}, \text{d})^{10}\text{Be}$   $Q_m = 0.558$

Not reported.

13.  $^9\text{Be}(\alpha, ^3\text{He})^{10}\text{Be}$   $Q_m = -13.763$

Not reported.

14.  $^{10}\text{B}(\text{n}, \text{p})^{10}\text{Be}$   $Q_m = 0.228$

See (1948EG1A, 1955JA18, 1964PA1M).

15.  $^{10}\text{B}(\text{t}, ^3\text{He})^{10}\text{Be}$   $Q_m = -0.536$

Table 10.5: Radiative transitions in  ${}^9\text{Be}(\text{d}, \text{p}){}^{10}\text{Be}$ <sup>a</sup>

Transition	$\Delta J^\pi$	$E_\gamma$ <sup>b</sup> (keV)	Mult.	Branch (%)	$\tau_m$ (psec)
$3.37 \rightarrow \text{g.s.}$	$2^+ \rightarrow 0^+$	$3374 \pm 10$	E2	100	$0.15 \pm 0.03$
$5.96 \rightarrow \text{g.s.}$	$1^- \rightarrow 0^+$	$5965 \pm 10$	E1	$48 \pm 2$ <sup>c</sup>	$< 0.08$
$5.96 \rightarrow 3.37$	$1^- \rightarrow 2^+$	$2584 \pm 15$		$52 \pm 2$ <sup>c</sup>	
$6.18 \rightarrow \text{g.s.}$	$0^+ \rightarrow 0^+$		E0	(100)	$> 0.5$
$6.18 \rightarrow 3.37$	$0^+ \rightarrow 2^+$				
$6.26 \rightarrow \text{g.s.}$	$2^- \rightarrow 0^+$			$< 0.4$ <sup>d</sup>	
$6.26 \rightarrow 3.37$	$2^- \rightarrow 2^+$			$> 99.6$ <sup>d</sup>	

<sup>a</sup> (1963WA03, 1963WA17, 1964WA05, 1965WA1P).

<sup>b</sup> Corrected for Doppler shift.

<sup>c</sup>  $(22 \pm 6)/(78 \pm 12)$  (1958ME81),  $(41 \pm 4)/(59 \pm 4)$  (1965FO1G); however, see  ${}^9\text{Be}(\text{n}, \gamma)$ .

<sup>d</sup> See also (1965FO1G).

Not reported.

$$16. {}^{11}\text{B}(\text{n}, \text{d}){}^{10}\text{Be} \quad Q_m = -9.004$$

Not reported.

$$17. {}^{11}\text{B}(\text{p}, 2\text{p}){}^{10}\text{Be} \quad Q_m = -11.228$$

In the summed proton spectrum, structure is observed corresponding to  $Q = -10.9 \pm 0.35$ ,  $-14.7 \pm 0.4$ ,  $-21.1 \pm 0.4$ ,  $-35 \pm 1$  MeV (1958TY49, 1962GA09, 1962GA23, 1962GO1P, 1964TI02, 1965RI1A, 1966TY01). See also (1958MA1B, 1963BE42, 1963RI1B, 1964BA1C), and  ${}^{11}\text{B}$ .

$$18. {}^{11}\text{B}(\text{d}, {}^3\text{He}){}^{10}\text{Be} \quad Q_m = -5.735$$

Not reported.



See ([1961HO01](#)).



See ([1963SA1K](#)).



See  $^{14}\text{C}$ .

**$^{10}\mathbf{B}$**   
(Figs. 20 and 22)

GENERAL: See (1959BA1F, 1959BR1E, 1960TA1C, 1961TR1B, 1962IN02, 1963BU1C, 1963KU03, 1963ME01, 1963MO1F, 1963OL1B, 1963VL1A, 1963WA1M, 1964AM1D, 1964BA29, 1964FR1D, 1964GR1J, 1964MA1G, 1964NE1E, 1964OL1A, 1964ST1B, 1964VA1D, 1965FA1C, 1965NE1C).

*Ground State:*

$$\mu = +1.8007 \text{ nm (1965FU1G).}$$

$$Q = +0.08 \text{ b (1965FU1G).}$$



Six resonances are observed in the range  $E_\alpha = 0.5$  to  $2.6$  MeV, corresponding to  ${}^{10}\text{B}^*(4.76 - 6.06 \text{ MeV})$ : see Table 10.8. No other resonances appear for  $E_\alpha < 3.8$  MeV ( ${}^{10}\text{B}^*(6.74)$ ) (1957ME27, 1961SP02).

The 4.77 MeV state decays mainly to  ${}^{10}\text{B}^*(0.7)$ : the ground state decay is  $< 3\%$  (1957WA07), 8% (1957ME27): see also Table 10.7. The angular distribution of  $\gamma$ -rays indicates  $J = 2^+$ , with  $E2/M1 = 1.8$  (1957ME27), 0.64 (1957WA07). The measured  $\omega\Gamma_s \equiv \omega\Gamma_\gamma\Gamma_\alpha/(\Gamma_\gamma + \Gamma_\alpha)$  is 0.05 eV:  $\Gamma(M1) \approx 0.01\Gamma_w$ ,  $\Gamma(E2) \gtrsim 10\Gamma_w$ , consistent with the  $\Delta T$  selection rule for M1 and considerable collective enhancement for E2 (1957ME27, 1957WA07, 1958ME81, 1963WA17).

The angular distribution of the  $\gamma$ -rays from the 5.11-MeV state can be made consistent with  $J = 2^-$ ;  $T = 0$  if  $M2/E1 \approx 0.01$  is assumed (1957ME27, 1958ME81). The distributions from  $E_x = 5.17$  MeV are consistent with  $J^\pi = 2^+$  (1957ME27): observations in  ${}^9\text{Be}(d, n){}^{10}\text{B}$  indicate  $\Gamma_\alpha \approx \Gamma_\gamma = 1.3$  eV and hence  $T = 1$  (1958ME81, 1959WA16, 1962WA21).

A study of  $\alpha$ -capture near  $E_\alpha = 1.18$  MeV shows the formation of a broad  $J = 1$ ;  $T = 0$  state ( $E_x = 5.18$  MeV,  $\Gamma_{c.m.} = 200 \pm 30$  keV), and its subsequent decay via 3.44 MeV  $\gamma$ -rays to the  $J^\pi = 0^+$ ;  $T = 1$  state at 1.74 MeV. The observed width corresponds to 0.86 of the single particle limit for s-wave  $\alpha$ -formation (1961SP02: see, however, (1962DE10)). See also (1962WA21).

For the 5.92 MeV level,  $J^\pi = 2^\pm$ ,  $3^+$ , and  $4^+$  are possible. Only  $J^\pi = 4^+$  gives a satisfactory account of the angular distribution from the 6.03 MeV level (1957ME27).

- |  |                |               |
|--|----------------|---------------|
| 2. (a) ${}^6\text{Li}(\alpha, p){}^9\text{Be}$   | $Q_m = -2.126$ | $E_b = 4.461$ |
| (b) ${}^6\text{Li}(\alpha, d){}^8\text{Be}$      | $Q_m = -1.567$ |               |
| (c) ${}^6\text{Li}(\alpha, \alpha){}^6\text{Li}$ |                |               |
| (d) ${}^6\text{Li}(\alpha, n){}^9\text{B}$       | $Q_m = -3.977$ |               |

Table 10.6: Energy levels of  $^{10}\text{B}$ 

$E_x$ (MeV $\pm$ keV)	$J^\pi; T$	$\Gamma$ (keV) or $\tau_m$ (psec)	Decay	Reactions
g.s.	$3^+; 0$	stable	—	1, 3, 4, 5, 6, 11, 12, 13, 14, 15, 20, 21, 22, 27, 28, 30, 33, 35
$0.7173 \pm 0.8$	$1^+; 0$	$1010 \pm 20$ psec	$\gamma$	1, 3, 4, 5, 6, 11, 12, 19, 20, 22, 26, 28, 30, 33, 35
$1.740 \pm 2$	$0^+; 1$	$0.15 \pm 0.02$ psec	$\gamma$	1, 3, 6, 11, 12, 20, 26, 28, 30, 33, 35
$2.154 \pm 3$	$1^+; 0$	$1.35 \pm 0.16$ psec	$\gamma$	1, 3, 6, 11, 12, 20, 22, 28, 30, 33, 35
$3.585 \pm 4$	$2^+; 0$	$0.096 \pm 0.036$ psec	$\gamma$	4, 6, 11, 12, 20, 22, 28, 30, 33
$4.774 \pm 3$	$(2^+); 0$	$< 10$	$\gamma, \alpha$	1, 4, 5, 11, 12, 20, 22, 30, 33
$5.114 \pm 4$	$(2^-); 0$	1.2	$\gamma, \alpha$	1, 11, 12, 20, 22, 30, 33
$5.166 \pm 4$	$2^+; 1$	0.003	$\gamma, \alpha$	1, 11, 12, 20, 30
$5.183 \pm 8$	$1^+; 0$	$120 \pm 20$	$\gamma, \alpha$	1, 2, 6, 11, 20, 22, 33
$5.923 \pm 4$	$2^+; 0$	$< 5$	$\gamma, \alpha$	1, 2, 11, 12, 20, 22, 30, 33
$6.029 \pm 4$	$4^+$	$< 1$	$\gamma, \alpha$	1, 2, 11, 12, 18, 20, 22, 30, 33
$6.133 \pm 4$		$< 5$	$\alpha$	2, 11, 12, 20, 22, 30
$6.566 \pm 6$		$\approx 70$	$\alpha$	2, 11, 12, 20, 22, 30
6.884	$1^-; 0$	$\approx 140$	$\gamma, p, d, \alpha$	2, 6, 8, 10
7.00		$95 \pm 10$	$p, d, \alpha$	2, 10, 20, 22
$7.431 \pm 10$	$2^-; 0$	$140 \pm 30$	$\gamma, p, d, \alpha$	6, 10
$(7.468 \pm 10)$	$(2^+)$	80	$\gamma, p$	6, 8, 18
$7.479 \pm 2$	$(2^-; 1)$	$72 \pm 4$	$\gamma, p$	6, 8, 20, 30
$7.561 \pm 1$	$0^+; 1$	$3.3 \pm 0.3$	$\gamma, p$	6, 8, 30
$7.62 \pm 50$	$(1^+; 0)$	$220 \pm 50$	$p, d, \alpha$	8, 10
$7.77 \pm 30$	$2^-; 1$	$210 \pm 60$	$\gamma, p, \alpha$	2, 6, 8
$8.07 \pm 100$	$(2^-; 0)$	$800 \pm 200$	$p, d, \alpha$	10
$8.892 \pm 6$	$3^{(-)}; (1)$	$84 \pm 7$	$n, p$	7, 8, 16
$8.896 \pm 2$	$2^+; 1$	$36 \pm 2$	$\gamma, p, \alpha$	6, 8, 10
9.7	; 1	$\approx 600$	$n, p, \alpha$	7, 10
$10.83 \pm 30$	; 1	$\approx 500$	$\gamma, n, p, \alpha$	6, 7, 10, 16
(11.4)	(+)		$\gamma$	16, 18, 20, 28
(14)	(+)		$\gamma$	16, 18, 20
18.6		$< 500$	$\gamma, {}^3\text{He}$	4
18.8	$(2^+, 1^-; 1)$	$< 600$	$\gamma, {}^3\text{He}, \alpha$	4
19.3	$(2^+, 1^-; 1)$	350	$\gamma, n, {}^3\text{He}, \alpha$	4
20.1		broad	$n, p, {}^3\text{He}$	4

Table 10.7: Electromagnetic transitions in  $^{10}\text{B}$ 

Initial State	$\Gamma_\gamma$ (total) (eV)	Relative intensities to final states at:						References
		g.s.	0.72	1.74	2.15	3.59	5.18	
0.72	$6.6 \times 10^{-7}$	100						Table 10.20 (1964SI03, 1965LO04)
1.74	$4.4 \times 10^{-3}$	< 2	100					(1961SP04, 1964HO02, 1965WA1P)
2.15	$4.9 \times 10^{-4}$	22	27	51				(1958ME81, 1964BR06, 1964HO02,
3.59	$6.9 \times 10^{-3}$	15	65	< 3	20			1964SI03, 1965WA1P)
4.77	0.03	8	92					Table 10.8, (1963WA17)
5.11	0.07	96	4					Table 10.8, (1963WA17)
5.17	0.6	5.5	29.5		65			Tables 10.8 and 10.19, (1963WA17)
5.18	0.06			100				Table 10.8
5.92		100						Table 10.8
6.03		100						Table 10.8
6.4	0.75 <sup>c</sup>							(1964JA03)
6.88	(4.8)	< 3	24	56	17	4	< 11	Table 10.13
7.43	2.4	< 2	1.3	< 0.1	0.62	0.5	< 1	Table 10.13 <sup>d</sup> (1965SP04)
7.47	11							
7.48	25.8	25	0.3	< 0.1	0.49	0	< 1	Table 10.13 <sup>d</sup>
7.56	8.5	< 0.2	6.7	< 0.3	0.8	< 0.2	1.0	Table 10.13 <sup>d</sup>
7.77	8.5	6.6	0.9	< 0.1	0.3	0.3	0.4	Table 10.13 <sup>d</sup>
8.89	0.6							(1964GR40)
11.8 <sup>a</sup>								
14 <sup>a</sup>								
18.6 <sup>b</sup>								
18.8 <sup>b</sup>								
19.3 <sup>b</sup>								

<sup>a</sup> See (1962ED02).

<sup>b</sup> See Table 10.10.

<sup>c</sup>  $B(\text{E}2\uparrow) = 29 \text{ fm}^4$ .

<sup>d</sup> Partial  $\Gamma_\gamma$  in eV (1964HO02).

Table 10.8: Levels of  $^{10}\text{B}$  from  $^6\text{Li}(\alpha, \gamma)^{10}\text{B}$  ([1957ME27](#), [1961SP02](#))

$E_{\text{res}}$ (keV)	$E_x$ (MeV)	$J^\pi; T$	$\Gamma_{\text{lab}}$ (keV)	$E_\gamma$ <sup>a</sup> (MeV)	Branch %	$\omega\Gamma_s$ <sup>b</sup> (eV)	
500 $\pm$ 25 <sup>c</sup>	4.76	(2 <sup>+</sup> ); 0		4.76	8		
				4.05	92	0.05 <sup>d</sup>	E2/M1 = 1.8
1085	5.112	2 <sup>-</sup> ; 0	2 <sup>e</sup>	5.1	96	0.10	M2/E1 $\approx$ 0.01
				4.4	4	0.005	
1175	5.167	2 <sup>+</sup> ; 1	< 0.5 <sup>e</sup>	5.16	7	0.04	E2/M1 = 0.02
				4.44	29	0.15	E2/M1 = 0.01
				3.01	64	0.32	E2/M1 = 0.01
1210 $\pm$ 35	5.188	1 <sup>+</sup> ; 0	340 $\pm$ 50	3.44	$\approx$ 100	f	
2435	5.923	2 <sup>+</sup>	20	5.9	100		E2/M1 = 0.01 <sup>g</sup>
2605	6.025	4 <sup>+</sup>	< 1.5 <sup>e</sup>	6.0	100		E2/M1 = 9.0

<sup>a</sup> Primary radiation: see Table [10.7](#).

<sup>b</sup>  $\omega\Gamma_s \equiv \omega\Gamma_\alpha\Gamma_\gamma/(\Gamma_\alpha + \Gamma_\gamma)$ .

<sup>c</sup> ([1953WI32](#), [1954JO09](#)).

<sup>d</sup> ([1957WA07](#)): E2/M1 = 0.64,  $\gamma$ (g.s.) < 3%;  $\Gamma_\gamma/\Gamma < 4 \times 10^{-4}$  ([1965RO01](#)).

<sup>e</sup> See ([1958ME81](#)) and  $^9\text{Be}(\text{d}, \text{n})^{10}\text{B}$ .

<sup>f</sup>  $\Gamma_\gamma = 0.06 \pm 0.03$  ([1961SP02](#)).

<sup>g</sup> 0.01 or 10 ([1957ME27](#)).

Reported anomalies in the elastic scattering are listed in Table [10.9](#) ([1962BA03](#), [1962DE10](#), [1965SI1B](#)). Angular distributions indicate  $J^\pi = 1^+$  and  $2^+$ , respectively for the 5.18 and 5.92 MeV states. The  $1^+$  assignment supports the proposal by ([1961TR1B](#)) that the 5.18 MeV state is a member of the doublet formed by two-nucleon excitation into the 2s shell, whose other member is the  $0^+; T = 1$  state at 7.56 MeV.

The excitation functions for  $\alpha_0$  and  $\alpha_1$  (to the 2.18 MeV state of  $^6\text{Li}$ ) particles ([1963BL20](#):  $E_\alpha = 9.50$  to  $12.50$  MeV),  $d_0$  particles (to the ground state of  $^8\text{Be}$ ) ([1963BL20](#):  $E_\alpha = 9.50$  to  $11.4$  MeV) and neutrons ([1963ME08](#): from threshold,  $E_\alpha = 6.623$  MeV, to  $15.5$  MeV) do not show resonance structure. See also ([1956WA29](#), [1964DE1K](#)).

- 3. (a)  $^6\text{Li}(^6\text{Li}, \text{d})^{10}\text{B}$   $Q_m = 2.989$
- (b)  $^6\text{Li}(^7\text{Li}, \text{t})^{10}\text{B}$   $Q_m = 1.994$
- (c)  $^6\text{Li}(^9\text{Be}, ^5\text{He})^{10}\text{B}$   $Q_m = 1.933$

Table 10.9:  $^{10}\text{B}$  levels from  $^6\text{Li}(\alpha, \alpha)^6\text{Li}$ 

$E_\alpha$ (MeV)	$E_\alpha$ (MeV)	$E_x$ (MeV)	$\Gamma_{\text{lab}}$ (keV)	$J^\pi; T$	$\theta_\alpha^2$	References
1.21		5.19	175	$1^+; 0$	$\approx 2$	(1962DE10)
2.44	2.44	5.93	$\approx 30$	$2^+; 0$	$\approx 0.05$	(1962DE10, 1965SI1B)
	2.605	6.03				(1965SI1B)
	2.80	6.14				(1965SI1B)
3.52	3.51	6.58	120			(1962BA03, 1965SI1B)
	4.05	6.89				(1965SI1B)
4.18	4.24	7.01	230			(1962BA03, 1965SI1B)
	5.50	7.76				(1965SI1B)

Reactions (a) and (b) have been studied up to 6 MeV bombarding energy. Deuteron and triton groups have been observed leading to the first four states of  $^{10}\text{B}$  with intensities which depend on the incident energy: typically the group to the 0.7 MeV state is very strong and that to the 1.74 MeV state is very weak (1960MO17, 1961BR35, 1961MO02, 1964KI02). See also (1962BE24, 1962BU1C, 1962MC12, 1963LE09, 1964BL1C). For reaction (c) see (1962MC12, 1963NO02).

4. (a) $^7\text{Li}(^3\text{He}, \text{n})^9\text{B}$	$Q_m = 9.349$	$E_b = 17.786$
(b) $^7\text{Li}(^3\text{He}, \text{p})^9\text{Be}$	$Q_m = 11.199$	
(c) $^7\text{Li}(^3\text{He}, \text{d})^8\text{Be}$	$Q_m = 11.759$	
(d) $^7\text{Li}(^3\text{He}, \alpha)^6\text{Li}$	$Q_m = 13.325$	
(e) $^7\text{Li}(^3\text{He}, \gamma)^{10}\text{B}$	$Q_m = 17.786$	

The excitation curve for reaction (a) is smooth up to  $E(^3\text{He}) = 1.8$  MeV (1962SE1A, 1963DU12), and shows resonance behavior at  $E(^3\text{He}) = 2.2$  and 3.25 MeV: the 2.2 MeV resonance has  $\Gamma \approx 280$  keV; the 3 MeV resonance is broader (1963DI01, 1963DU12, 1964DI1C).

Capture  $\gamma$ -rays have been observed for  $E(^3\text{He}) = 0.8$  to 3.0 MeV. The excitation functions for the transitions to the ground and 4.77 MeV states show peaks at  $E(^3\text{He}) = 1.1$  and 2.2 MeV; those to the 0.72 and 3.59 MeV states show a broad maximum at 1.4 MeV. The observed gamma widths are comparatively large (see Table 10.10) (1965PA02).

The yield of protons (reaction (b)) is relatively flat for  $E(^3\text{He}) = 2.5$  to 4.8 MeV, with some indication of a weak maximum at  $\approx 3.3$  MeV at  $20^\circ$  and  $50^\circ$  (1961WO05). The yield of ground state  $\alpha$ -particles at  $8^\circ$  (reaction (d)) shows a broad maximum at  $\approx 2$  MeV, a minimum at 3 MeV, followed by a steep rise which flattens off between  $E(^3\text{He}) = 4.5$  and 5.5 MeV. Integrated  $\alpha_0$  and  $\alpha_1$  yields rise monotonically to 4 MeV and then tend to decrease. Angular distributions of  $\alpha_1$  in the

Table 10.10: Resonances in  ${}^7\text{Li} + {}^3\text{He}$ 

$E_{\text{res}}$ (MeV)	$\Gamma_{\text{c.m.}}$ (keV)	$E_x$ (MeV)	$\omega\Gamma^{\text{a}}$ for transition to				$\Gamma_\alpha$	References
			g.s.	0.72	3.59	4.77		
1.1	< 500	18.6	35			140		(1965PA02)
1.4	< 600	18.8		100	100		< 80	(1965PA02, 1965PA03)
2.2	280 – 420	19.3	85			200	< 20	b
3.25	broad	20.1						(1963DI01, 1964DI1C)

<sup>a</sup> Lower limits of  $(2J + 1)\Gamma_\gamma$  in eV.

b (1963DI01, 1963DU12, 1964DI1C, 1965PA02, 1965PA03).

 Table 10.11: Slow neutron thresholds <sup>a</sup> in  ${}^7\text{Li}(\alpha, \text{n}){}^{10}\text{B}$ 

$E_\alpha(\text{threshold})$ (MeV)		$E_x$ (MeV)
(1957BI84)	(1963ME08)	
4.379 ± 0.006	4.380 ± 0.020	0
5.51	5.5 (11.90) (14.53)	0.72 4.77 6.42

<sup>a</sup> See also (1956RO06).

Table 10.12: Resonances in  ${}^9\text{Be}(\text{p}, \gamma){}^{10}\text{B}$ 

$E_{\text{p}}$ (MeV $\pm$ keV)	$E_{\text{x}}$ (MeV)	$\Gamma_{\text{c.m.}}$ (keV)	$J^\pi; T$	$\Gamma_{\text{p}}/\Gamma$	$\Gamma_\gamma^{\text{a}}$ (eV)	References
0.330	6.88	$\approx 140$	$1^-; 0$	0.30	4.8	(1956CL69)
$0.938 \pm 10$ (0.98)	7.43 (7.47)	$140 \pm 30$	$(2^-); 0$ $(2^+)$	0.7	2.4	(1962EL06, 1964HO02) (1962EL06)
$0.992 \pm 2$	7.48	$72 \pm 4$	$2^-; 1$	$\approx 0.65$	25.8	(1964HO02)
$1.0832 \pm 0.4$	7.562	$3.25 \pm 0.33$	$0^+; 1$	1.0	8.5	(1964BO13, 1964HO02)
1.29	7.75	$210 \pm 60$	$2^-; (1)$	$\approx 0.65$	8.5	(1964HO02)
$2.567 \pm 2$	8.896	$36 \pm 2$	$2^+; 1$			(1953MA1A, 1956MA55)
4.72	10.83	$\approx 500$				(1952HA10)

<sup>a</sup> See Table 10.13 for decay schemes.

range  $E({}^3\text{He}) = 2$  to 5.5 MeV suggest that the reaction proceeds mainly by pickup (1965FO07: se also (1961WO05)). Angular distributions for  $E({}^3\text{He}) = 0.8$  to 3.0 MeV give indications of the resonances at  $E({}^3\text{He}) = 1.4$  and 2.2 MeV seen in  ${}^7\text{Li}({}^3\text{He}, \gamma)$ :  $J^\pi = 2^+$  or  $1^-$ ;  $T = (1)$  for both:  $\Gamma_\alpha$  is small (1965PA03).

$$5. {}^7\text{Li}(\alpha, \text{n}){}^{10}\text{B} \quad Q_{\text{m}} = -2.792$$

Observed slow neutron thresholds are listed in Table 10.11 (1957BI84, 1963ME08). At  $E_\alpha = 13.5$  and 13.9 MeV, angular distributions of the ground-state neutrons have been determined (1962KJ05). See also (1959HE1B, 1962GA1L).

$$6. {}^9\text{Be}(\text{p}, \gamma){}^{10}\text{B} \quad Q_{\text{m}} = 6.587$$

Parameters of observed resonances are listed in Tables 10.12 and 10.13. Table 10.7 summarizes the  $\gamma$ -transitions from this and other reactions.

The  $E_{\text{p}} = 0.33$  MeV resonance ( ${}^{10}\text{B}^*(6.88)$ ) is ascribed to s-wave protons because of its comparatively large proton width [see  ${}^9\text{Be}(\text{p}, \text{p})$ ] and because of the isotropy of the  $\gamma$ -radiation. The strong transition to  ${}^{10}\text{B}^*(1.74)$  requires E1 and hence  $J^\pi = 1^-$ ;  $T = 0$ .  $T = 0$  is also indicated by the large deuteron width. On the other hand, the strength of E1 transitions to  ${}^{10}\text{B}^*(0.7, 2.1)$  indicate a  $T = 1$  admixture of 20% or more (1956WI16, 1959ME85). A small  $P_2$  term in

Table 10.13: Radiative transitions in  ${}^9\text{Be}(\text{p}, \gamma){}^{10}\text{B}$ 

Initial state (MeV)	$\Gamma_\gamma(\text{tot})$ (eV)	Relative intensities to final states						Remarks	Ref.
		ground $3^+; 0$	0.72 $1^+; 0$	1.74 $0^+; 1$	2.15 $1^+; 0$	3.59 $2^+; 0$	5.18 $1^+; 0$		
6.88 $E_{\text{p}} = 0.33$	[10]	2.4	6	15	3	5	3.6	$ M ^2\text{E1}$	(1957BI75)
		15	40	100	45				(1955CA25)
		4.8	1.0	2.6	0.5		0.65		(1956CL69)
		< 0.7	6.4	15	4.5	$\approx 1$	< 2.9		(1959ME85)
			0.006	0.025	0.009	$\approx 0.005$			(1959ME85)
7.43 $E_{\text{p}} = 0.94$	[2.4]		(res)					$ M ^2\text{E1}$	(1962EL06)
		< 2	1.3	[< 0.14]	0.62	0.5	[< 1]		(1964HO02)
			0.013		0.013	0.03			(1964HO02)
7.48 $E_{\text{p}} = 0.99$	[25.8]	400	< 10	19	22	< 7	< 13	$ M ^2\text{E1}$	(1959ME85)
		25	0.3	[< 0.14]	0.49	0	[< 1]		(1964HO02)
		0.19	0.003		0.10				(1964HO02)
7.56 $E_{\text{p}} = 1.08$	4.8	1.0	2.6	0.5		0.65		$ M ^2\text{M1}$	(1956CL69)
		< 3	100	< 8	< 8	< 4	23		(1959ME85)
			76		9		15		(1961SP04)
	6.6 [8.5]	100	< 2	10		40			(1962ME1A)
		< 0.2	6.7	< 0.3	0.8	< 0.2	1.0		(1964HO02)
			1.0		0.24		3.5		(1964HO02)
7.75 $E_{\text{p}} = 1.29$	14.6 [8.5]	83%						$ M ^2\text{E1}$	(1963FU11)
		6.6	0.9	< 0.08	0.3	0.3	0.4		(1964HO02)
		0.044	0.008		0.006	0.013	0.08		(1964HO02)

the angular distribution at resonance suggests a d-wave admixture. The angular correlation of  $6.88 \rightarrow 0.72 \rightarrow \text{g.s.}$  is consistent with  $J^\pi = 1^-$  but does not exclude  $2^-$  (1964BI18).

The proton capture data near  $E_{\text{p}} = 1$  MeV appears to require at least 5 resonant states, at  $E_{\text{p}} = 938, (980), 992, 1086$  and  $1290$  keV. The narrow  $E_{\text{p}} = 1086$  keV level ( ${}^{10}\text{B}^*(7.56)$ ) is formed by p-wave protons,  $J^\pi = 0^+$  [see  ${}^9\text{Be}(\text{p}, \text{p})$  and  ${}^9\text{Be}(\text{p}, \alpha)$ ]. The isotropy of the gamma rays supports this assignment (1961TA02). The strong M1 transitions to  $J = 1^+; T = 0$  levels at  $0.7, 2.15$  and  $5.18$  MeV (Table 10.13) indicate  $T = 1$  (1959WA16).

The excitation function for ground-state radiation shows resonance at  $E_{\text{p}} = 992$  ( $\Gamma = 80$  keV) and  $1290$  keV ( $\Gamma = 230$  keV) (1962EL06, 1964HO02). Elastic scattering studies indicate s-wave formation and  $J = 2^-$  for both (1956MO90). For the lower level ( $E_x = 7.48$  MeV) the intensity of the g.s. capture radiation,  $\Gamma_\gamma = 25$  eV (1964HO02) indicates E1 and  $T = 1$ . The angular distribution of  $\gamma$ -rays,  $1 + 0.1 \sin^2 \theta$ , is consistent with s-wave formation with some d-wave admixture (1953PA22) or with some contribution from a nearby p-wave resonance (1956MO90); possibly a  $J^\pi = 2^+$  level at  $E_{\text{p}} = 980$  keV (1956MO90, 1962EL06: see, however, (1964HO02)).

The angular correlation of internal pairs is consistent with an E1/M1 mixture of 3 : 2 ([1962EL06](#)). Earlier difficulties with the  $T = 1$  assignment may be resolved if the (p, d) and (p,  $\alpha$ ) resonances are ascribed to another level ([1964HO02](#)).

The angular distribution of ground-state radiation at  $E_p = 1330$  keV is isotropic and  $\Gamma_\gamma = 14.6 \pm 1.5$  eV ([1963FU11](#)), 8.5 eV ([1964HO02](#)), supporting E1,  $T = 1$  for this level,  $E_x = 7.75$  MeV.

Transitions to  $^{10}\text{B}^*(0.7)$  show resonance at  $E_p = 992$ , 1290 and 938 keV,  $\Gamma = 155$  keV ([1962EL06](#), [1964HO02](#)). The latter is presumably also a resonance for (p, d) and (p,  $\alpha$ ). An assignment of  $J^\pi = 2^-$ :  $T = 0$  is consistent with the data, although the E1 radiation then seems somewhat too strong for a  $\Delta T = 0$  transition ([1964HO02](#)).

A resonance for capture radiation at  $E_p = 2.567 \pm 0.003$  ( $E_x = 8.896$  MeV) has a width of  $40 \pm 2$  keV and decays mainly via  $^{10}\text{B}^*(0.7)$  ([1953MA1A](#)). It appears from the width that this resonance corresponds to that observed in  $^9\text{Be}(p, \alpha)$ ,  $J = 2^+$ ;  $T = 1$  and not to the  $^9\text{Be}(p, n)$  resonance at the same energy ([1956MA55](#)). A fourth resonance is reported at  $E_p = 4.72 \pm 0.01$  MeV,  $\Gamma \approx 0.5$  MeV ([1952HA10](#)).

For the mean life of the 0.7 MeV state, see Table [10.20](#).

See also ([1959KA69](#), [1959SI1C](#), [1960GO23](#), [1960SI04](#), [1961RI08](#), [1962BL10](#), [1962HU05](#), [1963CO1K](#), [1964SI03](#)).

$$7. \ ^9\text{Be}(p, n)^9\text{B} \quad Q_m = -1.851 \quad E_b = 6.587$$

Resonances in neutron yield occur at  $E_p = 2.56$  and 4.6 MeV: see Table [10.14](#). There is some indication of a broad maximum near  $E_p = 3.5$  MeV; a peak reported at  $E_p = 4.9$  MeV for  $n_1$  neutrons may reflect the effect of this level ([1959MA20](#)). A sharp break at  $E_p = 6.55 \pm 0.03$  MeV is ascribed to a level in  $^9\text{B}$  at 4.04 MeV ([1964BA16](#)). Angular distributions in the range  $E_p = 2$  to 11 MeV are reported by ([1956MA55](#), [1961AL07](#), [1963KE03](#), [1965WA04](#)). Polarization studies have been made by ([1961CR1A](#), [1963KE03](#), [1965WA04](#)).

The  $E_p = 2.56$  MeV resonance is considerably broader than that observed at the same energy in  $^9\text{Be}(p, \alpha)$  and  $^9\text{Be}(p, \gamma)$  and the two resonances are believed to be distinct ([1956MA55](#)). The shape of the resonance and the magnitude of the cross section can be accounted for with  $J = 3^-$  or  $3^+$ : the former assignment is in better accord with charge symmetry and indicates correspondence with  $^{10}\text{Be}^*(7.38)$ . For  $J^\pi = 3^-$ ,  $\theta_n^2 = 0.135$ ,  $\theta_p^2 = 0.115$  ( $R = 4.47$  fm). The  $J^\pi = 2^+$  level should contribute about 10% to the cross section at  $E_p = 2.56$  MeV ([1962AL1A](#)). See also ([1963HA1G](#), [1963VA1C](#)).

$$8. \ ^9\text{Be}(p, p)^9\text{Be} \quad E_b = 6.587$$

Elastic scattering has been studied for  $E_p = 0.2$  to 2.6 MeV by ([1956DE33](#), [1956MO90](#)): see Table [10.15](#). Below  $E_p = 0.7$  MeV, only s-waves are present, exhibiting resonance at  $E_p = 330$

Table 10.14: Resonances in  ${}^9\text{Be}(\text{p}, \text{n}){}^9\text{B}$ 

$E_{\text{res}}$ (MeV $\pm$ keV)	$\Gamma_{\text{lab}}$ (keV)	$E_x$ (MeV)	$J^\pi; T$	$\sigma_{\text{max}}$ (mb)	References
2.562 $\pm$ 6	85 $\pm$ 10	8.892	3 $^{(-)}$ ; (1)	160	(1956MA55, 1962AL1A)
	100 $\pm$ 10				(1959GI47)
3.5	$\approx$ 700	9.7		240	(1956MA55, 1959MA20)
					(1959GI47)
4.72 $\pm$ 10	500	10.83		470	(1952HA10)
4.68 $\pm$ 30					(1955MA84)
4.6					(1959GI47)
4.62 $\pm$ 30					(1959MA20)
4.94 $\pm$ 30	$\approx$ 100			510	(1955MA84)
4.85					(1959GI47)
4.82 $\pm$ 60				a	(1959MA20)

<sup>a</sup> Resonance for n<sub>1</sub> only: possibly due to 3.5 MeV resonance.

keV,  $J = 1^-$  or  $2^-$ : both proton and deuteron widths are large, while  $\theta_\alpha^2$  is small. Further s-wave resonances, with  $J = 2^-$ , appear at  $E_p = 998$  and 1330 keV and a sharp p-wave resonance,  $J = 0^+$ , occurs at 1084 keV. The behavior of p-wave phase shifts indicates an additional  $J = 2^+$  resonance at 980 keV (1956MO90) or near 1100 keV (1956DE33). The behavior at  $E_p = 2.56$  MeV requires  $J \geq 2$  and a large proton width (1956DE33). See also (1959AJ76, 1963AN12, 1964HO02).

The yield of p<sub>0</sub> and p<sub>1</sub> (to 2.43 MeV state of  ${}^9\text{Be}$ ) has been determined by (1961RE03) in the range  $E_p = 3.6$  to 6.0 and 4.2 to 6.0 MeV, respectively, and for  $E_p = 5.7$  to 8 MeV by (1963BL20). Total cross sections and yields of the p<sub>1</sub> group have also been determined by (1964BI19) for  $E_p = 5$  to 15 MeV: a broad maximum near  $E_p = 8$  MeV is indicated. Total cross sections have also been measured at  $E_p = 10.2$ , 142 and 180 MeV (1961JO17, 1961TA06, 1962IG1A, 1963WI12, 1963WI1D). See also (1961JO18) and (1959AJ76).

Elastic scattering of polarized protons has been studied at  $E_p = 8.5$  and 11.4 MeV (1961RO05, 1961RO13); inelastic scattering has been studied at  $E_p = 7$  and 12 MeV (1965FU03). See also (1964BE03, 1964CR1B, 1965BO1L).

Table 10.15: Resonances in  ${}^9\text{Be}(\text{p}, \text{p}){}^9\text{Be}$ 

$E_{\text{res}}$ (keV)	$E_{\text{x}}$ (MeV)	$\Gamma_{\text{lab}}$ (keV)	$J^\pi$	$\Gamma_{\text{p}}/\Gamma$	$\theta_{\text{p}}^2$	References
330	6.88		$1^-$	0.30	0.5	(1956MO90)
$(980 \pm 10)^{\text{a}}$	(7.47)	90	$(2^+)$	0.9	0.07	(1956MO90)
998	7.48	$150 \pm 50$	$2^-$	0.65	0.02	(1956DE33, 1956MO90)
$1084 \pm 2$	7.56	3	$0^+$	1.0		(1956MO90)
$(1100)^{\text{a}}$	(7.58)	200		small		(1956DE33)
$1330 \left. \right\}$	7.79	$400 \pm 100$	$2^-$	0.65	0.05	(1956MO90)
$1350 \pm 50 \left. \right\}$		500	$2^-$	0.8		(1956DE33)
2560	8.89		$\geq 2$	large		(1956DE33)

<sup>a</sup> Postulated to account for p-wave structure near 1 MeV.

 Table 10.16: Resonances in  ${}^9\text{Be}(\text{p}, \text{d}){}^8\text{Be}$  and  ${}^9\text{Be}(\text{p}, \alpha){}^6\text{Li}$ 

$E_{\text{p}}$ (MeV)	$E_{\text{x}}$ (MeV)	$\Gamma_{\text{c.m.}}$ (keV)	$J^\pi; T$	$\Gamma_{\text{p}}/\Gamma$	$\theta_{\text{p}}^2$	$\theta_{\text{d}}^2$	$\theta_\alpha^2$	References
0.33	6.88		$1^-; 0$	0.30	0.5	0.4	0.05	(1956MO90)
(0.47)	(7.01)							(1949TH05, 1951NE03)
(0.68)	(7.20)							(1949TH05, 1951NE03)
0.94	7.43	140	$(2^-; 0)$	0.7	0.04	0.02		(1956WE37, 1964HO02)
1.15	7.62	$225 \pm 50$	$(1^+; 0)$	$\approx 0.4$	$\approx 0.1$			(1956WE37, 1964HO02)
1.65	8.07	$800 \pm 200$	$(2^-; 0)$	$\approx 0.07$	0.18	0.21		(1956WE37, 1964HO02)
(2.3)	(8.7)	$(\approx 220)$						(1956WE37, 1965MO27)
2.56	8.89	$36 \pm 2$	$2^+; 1$				<sup>a</sup>	(1956WE37)
3.5	9.7		$; 1$				<sup>a</sup>	(1959MA20)
4.49	10.63		$; 1$				<sup>a</sup>	(1959MA20)

<sup>a</sup> Resonance for  $\alpha_2$ , to  ${}^6\text{Li}^*(3.56), 0^+; T = 1$ .

Table 10.17: Levels from  $^{10}\text{B}$  of  $^9\text{Be}(\text{d}, \text{n})^{10}\text{B}$ 

$^{10}\text{B}^*$ (MeV)	$J^\pi$	$l_p$	$\theta^2$ <sup>a</sup>	References <sup>b</sup>
0	$3^+$	1	1	(1952AJ22, 1961MO15)
0.72	$1^+$	1	2.1	(1952AJ22, 1961MO15)
1.74	$0^+$	1	1.4	(1952AJ22, 1961MO15)
2.15	$1^+$	1	0.65	(1952AJ22, 1961MO15)
3.59	$2^+$	1	0.44	(1952AJ22, 1963FE1B)
4.77	$(2^+)$	(1)		(1963FE1B)
5.11	$(2^-)$	0		(1963RI08)
5.17	$(2^+)$	1		(1963RI08)
(5.18)				(1963FE1B)
5.58				(1952AJ22)
5.92		(1)		(1963FE1B)
6.02				(1963FE1B)
6.16				(1952AJ22, 1963FE1B)
6.38				(1952AJ22)
6.58				(1952AJ22)
(6.77)				(1952AJ22)

<sup>a</sup> Relative values (1952AJ22, 1960MA32).

<sup>b</sup> See also (1959AJ76).

$$9. \ ^9\text{Be}(\text{p}, \text{t})^7\text{Be} \quad Q_m = -12.079 \quad E_b = 6.587$$

See  $^7\text{Be}$ .

$$10. (\text{a}) \ ^9\text{Be}(\text{p}, \text{d})^8\text{Be} \quad Q_m = 0.559 \quad E_b = 6.587$$

$$(\text{b}) \ ^9\text{Be}(\text{p}, \alpha)^6\text{Li} \quad Q_m = 2.126$$

The (p, d) and (p,  $\alpha$ ) reactions have been studied in the range  $E_p = 0.8$  to 3.0 MeV by (1949TH05, 1951NE03, 1956MO90, 1956WE37),  $E_p = 1$  to 4.6 MeV by (1965MO27) and for  $E_p = 3.5$  to 12.5 MeV by (1963BL20); the (p,  $\alpha_2$ ) reaction, leading to  $^6\text{Li}^*(3.56)$ , has been studied from  $E_p = 2.3$  to 5.4 MeV by (1954MA26, 1956MA55, 1959MA20). Observed resonances are exhibited in Table 10.16.

Both alphas and deuterons are isotropic at the  $E_p = 0.33$  MeV resonance, confirming its s-wave formation: proton and deuteron widths are large, while  $\theta_\alpha^2$  is small ([1956MO90](#)). A strong maximum for  $\alpha$  and d appears at  $E_p = 0.93$  MeV,  $\Gamma = 130 \pm 30$  keV, followed by weaker maxima for d at  $E_p = 1.25, 1.65$  and  $2.3$  MeV. Alpha particles show a weak effect at the  $E_p = 2.56$  MeV,  $T = 1$  resonance, indicating a small isospin impurity ([1956WE37](#)). Angular distributions in the range  $E_p = 0.4$  to  $1.0$  MeV ([1951NE03](#)) and  $0.8$  to  $3.0$  MeV ([1956WE37](#)) show strong interference effects. Analysis of the latter data suggests contributions from three levels, at  $E_p = 0.938$  ( $2^-$ ),  $1.15$  ( $1^+$ ) and  $1.65$  MeV ( $2^-$ ) ([1964HO02](#)). There is no evidence of further structure in the yield of  $d_0$ ,  $\alpha_0$  or  $\alpha_1$  for  $E_p < 12$  MeV: in this range direct interaction appears to dominate ([1961RE03](#), [1963BL20](#), [1965MO27](#)). Polarization of  $d_0$  has been studied by ([1960BA26](#), [1961LA17](#), [1964BO33](#)).

The yield of  $3.56$  MeV  $\gamma$ -rays, associated with  $\alpha_2$  leading to  ${}^6\text{Li}^*(3.56)$  ( $J^\pi = 0^+; T = 1$ ), shows strong resonances at  $E_p = 2.56$  and  $4.49$  MeV and a broad rise at  $E_p = 3.5$  MeV, suggesting that these states, which are also observed in (p, n), have  $T = 1$  and are the analogues of  ${}^{10}\text{Be}^*(7.55, 9.26, 9.4)$  ([1959MA20](#)). See also ([1959LE27](#), [1961BE1E](#), [1961ST1D](#)).

$$11. {}^9\text{Be}(d, n){}^{10}\text{B} \quad Q_m = 4.363$$

Neutron groups are observed corresponding to  ${}^{10}\text{B}$  states listed in Table [10.17](#). There have been various reports of additional states: see ([1960HJ01](#), [1960JU04](#), [1961GA1G](#), [1962CO23](#), [1962MO12](#), [1963KO15](#)) and ([1959AJ76](#)). Thresholds for slow neutron production corresponding to  ${}^{10}\text{B}$  states from  $4.77$  to  $6.57$  MeV are reported in Table [10.18](#) ([1954BO79](#)). Angular distributions have been studied at many energies (see ([1959AJ76](#)) for a summary of the earlier work and ([1960BA46](#), [1961MO15](#), [1963FE1B](#), [1964BU14](#), [1965SI12](#))). The data, analyzed by stripping theories, show  $J \leq 3$  and even parity for the first five states of  ${}^{10}\text{B}$  ([1952AJ22](#), [1961MO15](#)) and for the  $5.17$  MeV state ([1962GA11](#), [1963RI08](#)) while the  $5.11$  MeV state is well fitted by  $l_p = 0$  ([1963RI08](#)).

Observed  $\gamma$ -transitions are listed in Tables [10.7](#) and [10.19](#). Reported values of the mean life of the  $0.72$  MeV state are given in Table [10.20](#). With an intermediate coupling parameter  $a/K = 4.75$ , a mean E2 lifetime of  $4$  nsec is predicted: the experimental value indicates either a lower value of  $a/K$  or some collective enhancement ([1957FR1B](#), [1957KU58](#), [1962LO02](#)). [Since the  ${}^{10}\text{C}$   $\beta$ -decay is allowed,  $J = 0^+, 1^+$  for  $E_x = 0.72$  MeV; the  $\gamma$ -transition from  $E_x = 1.74$  MeV established  $J = 1^+$ .]

The  $1.74$  MeV state,  $J = 0^+; T = 1$  analogue of  ${}^{10}\text{Be}$  and  ${}^{10}\text{C}$ , decays via the  $0.72$  MeV state. The  $2.15$  MeV state decays relatively strongly to  $E_x = 1.74$  MeV, arguing against  $J = 0, 2, 3$ : therefore  $J = 1^+$ . The E2 branch to the ground state is relatively strong compared to IPM predictions ([1957KU58](#)). The mean life is  $1.35 \pm 0.16$  psec ([1965WA1P](#)).

Correlation measurements in the cascade  $E_x = 3.59 \rightarrow 0.7 \rightarrow \text{g.s.}$  exclude  $J = 0$  or  $3$  for the  $3.59$  MeV state ([1956SH94](#)); a spin 1 assignment would permit a strong transition to  $E_x = 1.74$  MeV, therefore  $J = 2^+$ . The  $3.59 \rightarrow 0.7$  transition is either M1/E2 =  $93/7$  or pure E2 ([1956SH94](#)). The n- $\gamma$  correlation is isotropic ([1962GA11](#)). The large intensity relative to the

Table 10.18: Slow neutron thresholds in  ${}^9\text{Be}(\text{d}, \text{n}){}^{10}\text{B}$   
(1954BO79)

$E_{\text{d}}$ (MeV)	$E_{\text{x}}$ (MeV)	$\Gamma$ (keV)
0.52	4.79	< 10
0.92 <sup>a</sup>	5.12	< 10
0.99 <sup>a</sup>	5.17	< 10
1.92	5.93	< 10
2.08	6.06	< 10
2.20	6.16	< 20
2.53	6.43	
2.70	6.57	$\approx 30$

<sup>a</sup>  $0.921 \pm 0.009$  and  $0.989 \pm 0.009$  MeV (1962WA21).

Table 10.19: Gamma rays observed in  ${}^9\text{Be}(\text{d}, \text{n}){}^{10}\text{B}$

Transition	$E_{\gamma}$ (keV)	Branching ratio <sup>d</sup> (%)
$0.7 \rightarrow \text{g.s.}$	$716.6 \pm 1$ <sup>a</sup>	100
$1.7 \rightarrow \text{g.s.}$		< 2
$1.7 \rightarrow 0.7$	$1022 \pm 2$ <sup>a</sup>	100
$2.15 \rightarrow \text{g.s.}$	$2152 \pm 15$ <sup>a</sup>	22
$2.15 \rightarrow 0.7$	$1433 \pm 5$ <sup>a</sup>	27
$2.15 \rightarrow 1.7$	$413.5 \pm 1$ <sup>a</sup>	51
$3.59 \rightarrow \text{g.s.}$	$3583 \pm 13$ <sup>b,c</sup>	15 <sup>e</sup>
$3.59 \rightarrow 0.7$	$2872 \pm 15$ <sup>b,c</sup>	65 <sup>e</sup>
$3.59 \rightarrow 2.15$		20
$5.17 \rightarrow \text{g.s.}$	$5159 \pm 16$ <sup>b</sup>	$5.5 \pm 0.7$ <sup>b</sup>
$5.17 \rightarrow 0.7$	$4461 \pm 13$ <sup>b,c</sup>	$29.5 \pm 2$ <sup>b</sup>
$5.17 \rightarrow 2.15$	$3028 \pm 15$ <sup>b</sup>	$65 \pm 2$ <sup>b</sup>

<sup>a</sup> (1949RA02).

<sup>b</sup> (1963WA17): Doppler corrected.

<sup>c</sup> M1 + E2 (1964WA05).

<sup>d</sup> See Table 10.7.

<sup>e</sup> Ratio 1 :  $4.42 \pm 0.15$  (1964BR06).

Table 10.20: Mean life of  $^{10}\text{B}^*(0.72)$ 

$\tau_m$ (nsec)	Reaction	References
$0.7 \pm 0.2$	$^9\text{Be}(\text{d}, \text{n})^{10}\text{B}$ ibid ibid ibid and $(\text{p}, \text{p}')$ $^{10}\text{B}(\text{p}, \text{p}')^{10}\text{B}^*$ ibid ibid ibid ibid	(1953TH14)
$0.85 \pm 0.20$		(1956SE08)
$1.18 \pm 0.33$		(1961KN02) <sup>a</sup>
$0.96 \pm 0.10$		(1958DA11)
$1.04 \pm 0.02$		(1962LO02)
$1.05 \pm 0.10$		(1957BL02)
$0.90 \pm 0.1$		(1958HO97)
$0.94 \pm 0.05$		(1959BI10)
$1.00 \pm 0.03$		(1962GA15)
$1.0 \pm 0.1$		(1961SA21)
$1.013 \pm 0.015$		mean

<sup>a</sup> See also (1958GO47, 1958KN1B, 1960GO23, 1961RI08, 1962LO02).

ground state transition does not fit the IPM (1963WA17). The mean life is  $0.096 \pm 0.036$  psec (1965WA1P). See also (1959CH28, 1959GO78).

The 4.77 MeV state has  $J^\pi = 2^+$ ;  $T = 0$  [see  $^6\text{Li}(\alpha, \gamma)$ ]. No gamma radiation is observed in the present reaction:  $\Gamma_\gamma/\Gamma < 0.05$ . The relative weakness of the ground-state branch and the absence of the present level in the IPM scheme suggests a collective excitation based on  $E_x = 0.72$  MeV (1963WA17).

Three levels exist near 5 MeV, at 5.11, 5.17 and 5.18 MeV. For the  $E_x = 5.17$  MeV level, the small  $\Gamma_\alpha$  and large  $\Gamma_\gamma$  suggest  $T = 1$  (1959WA16) [confirmed in  $^{10}\text{B}(\text{d}, \text{d}')$ ]:  $\Gamma_\alpha \approx \Gamma_\gamma = 0.6$  eV (1963WA17),  $\Gamma_\gamma/\Gamma = 0.7 \pm 0.35$  (1962WA21),  $\Gamma_\gamma/\Gamma = 1 \pm 0.2$  (1963RI08). The mean life is  $< 0.08$  psec (1965WA1P). Angular correlations confirm the even parity assignment and indicate  $J^\pi = 2^+$  (1962GA11, 1963WA17). The 5.18 MeV level is excited only weakly, if at all, in the present reaction (1963FE1B, 1963RI08).

See also (1959HA1K, 1959LE1E, 1959SI1A, 1961HO1D, 1962LE1C, 1963MO1L, 1964BA1V, 1964SI02, 1965MA1K).

$$12. \ ^9\text{Be}(^3\text{He}, \text{d})^{10}\text{B} \quad Q_m = 1.094$$

At  $E(^3\text{He}) = 5.7, 8.8$  and  $10.2$  MeV, deuteron groups are observed corresponding to the ground state and to states at  $0.72, 1.74, 2.15$  and  $3.58$  MeV (1959HI69, 1960HI08: see Table 10.21).

Table 10.21: Levels of  $^{10}\text{B}$  from  $^9\text{Be}(^3\text{He}, \text{d})^{10}\text{B}$  ([1960HI08](#))<sup>a</sup>

$E_x$ (MeV $\pm$ keV)	$J^\pi$	$l_p$	$R$ (fm)	$\theta^2$ <sup>b</sup>
0	$3^+$	1	7.0	2.6
$0.717 \pm 10$	$1^+$	1	7.0	6.4
$1.744 \pm 10$	$0^+$	1	7.0	9.9
$2.156 \pm 10$	$1^+$	1	5.5	3.1
3.58	$2^+$	1	6.0	1.3

<sup>a</sup> See also ([1965CR1C](#), [1965YO1E](#)).

<sup>b</sup> Relative reduced widths,  $E(^3\text{He}) = 10.23$  MeV.

Angular distributions of deuterons to these states have been determined at a number of energies up to  $E(^3\text{He}) = 25$  MeV ([1959HI69](#), [1960HI08](#), [1960WE04](#)). Deuteron groups have also been observed to  $^{10}\text{B}^*(4.77, 5.11, 5.17, 5.92, 6.03, 6.13$  and  $6.57$  MeV). There is no indication of earlier reported levels at  $E_x = 5.58$  and  $6.40$  MeV ([1965CR1C](#), [1965YO1E](#)). See also ([1962WE1C](#)).

13.  $^9\text{Be}(\alpha, \text{t})^{10}\text{B}$   $Q_m = -13.227$

See ([1960GO04](#), [1960VL03](#), [1962WE1C](#)).

14.  $^9\text{Be}(^{14}\text{N}, ^{13}\text{C})^{10}\text{B}$   $Q_m = -0.963$

See ([1964BO1M](#)).

15.  $^{10}\text{Be}(\beta^-)^{10}\text{B}$   $Q_m = 0.555$

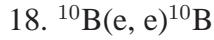
See  $^{10}\text{Be}$ .

16. (a)  $^{10}\text{B}(\gamma, \text{n})^9\text{B}$   $Q_m = -8.438$   
 (b)  $^{10}\text{B}(\gamma, \text{p})^9\text{Be}$   $Q_m = -6.587$   
 (c)  $^{10}\text{B}(\gamma, \text{d})^8\text{Be}$   $Q_m = -6.028$   
 (d)  $^{10}\text{B}(\gamma, \alpha)^6\text{Li}$   $Q_m = -4.461$

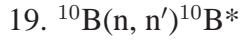
Energy levels are reported in  $^{10}\text{B}$  [from reaction (a)] at  $E_x = 10.25, 10.75, 11.85, 12.25, 14.7$  and  $16.9$  MeV, assuming that the neutron transitions are to the ground state of  $^9\text{B}$ . The giant dipole resonance appears to peak at  $11$  MeV ([1962FI07](#)). Cross sections have been measured with monoenergetic  $\gamma$ -rays for  $E_\gamma = 8.9$  to  $10.8$  MeV: the value at  $8.9$  MeV leads to  $\Gamma_\gamma = 0.6 \pm 0.3$  eV for the  $E_x = 8.89$  MeV level(s) ([1964GR40](#)). For reactions (b), (c), (d), see ([1959AJ76](#)) and ([1962CH26](#), [1962VO1D](#)).



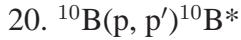
See ([1964LO1C](#)).



Elastic scattering at  $\theta = 180^\circ$  gives evidence of an M3 contribution ([1965GO1K](#): see also ([1965RA1D](#))). At  $E_e = 41.5$  MeV ( $\theta = 180^\circ$ ) evidence is reported for the M1 excitation of three states with  $J^\pi = 2^+, 3^+$  or  $4^+$  at  $E_x = 7.9, 11.8$  and  $14.0$  MeV, with  $\Gamma_\gamma$  in the range  $10$  to  $40$  eV ([1962ED02](#)). At  $E_e = 55$  MeV, transitions are observed to  $^{10}\text{B}^*(6.02 \pm 0.02$  and  $7.48 \pm 0.02$  MeV), the latter with  $\Gamma = 40$  keV,  $\Gamma_\gamma(\text{M1}) = 11 \pm 2$  eV, assuming  $J^\pi = 2^+$  ([1965SP04](#): see  $E_x = 7.47$ , Table 10.15). See also ([1959ME24](#), [1962BA1D](#), [1963GO04](#), [1963RO1M](#)). At  $100 - 200$  MeV, ([1965FR07](#)) find strong E2 excitation of the  $6.02$  MeV level.



See ([1956DA23](#), [1960AN14](#), [1963GL1F](#)).



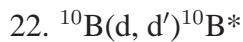
Excited states observed in inelastic scattering are listed in Table 10.22. Levels observed at  $E_x = 5.92, 6.03, 6.13$  and  $6.55$  MeV correspond well with those reported in  $^9\text{Be}(\text{d}, \text{n})^{10}\text{B}$ : no level at  $E_x = 6.43$  MeV is seen in the present reaction. A broad level at  $E_x = 7.00$  MeV may correspond to a peak reported in  $^9\text{Be}(\text{p}, \text{d})$  at  $E_p = 0.48$  MeV. The  $J = 2^-; T = 1, 7.48$  MeV state is seen here, but not in  $(\text{d}, \text{d}')$  ([1964AR04](#)). At  $E_p = 17.9$  MeV angular distributions are strongly peaked forward. Comparison of  $\sigma(\text{p}, \text{p}')$  with E2 transition rates suggests that the strongest  $(\text{p}, \text{p}')$  transitions correspond to states having large E2 coupling to the ground state: the strength of the  $6.04$  MeV  $4^+$  level is particularly significant ([1962SC12](#)),  $B(\text{E2}) = 29$  fm $^4$  ([1964JA03](#)). At  $E_p = 185$  MeV, the  $Q = -6.1$  MeV peak is dominant. Others are observed corresponding to

$E_x = 0, 2.2, 3.6, 5.2, 7.55, 11.0$  and  $13.0$  MeV. The angular distribution of the  $Q = -7.55$  MeV group is similar to that of the  $Q = -3.6$  MeV group of  ${}^6\text{Li}$  ([1965HA17](#)). Inelastic cross sections at  $E_p = 153$  MeV yield  $|M|^2 = 11, 0.02,$  and  $0.66$  for  $\gamma$ -transitions from  ${}^{10}\text{B}^*(2.15)$  to g.s.,  $0.7$  and  $1.74$  ([1961CL09](#)).

At  $E_p > 3$  MeV,  $\gamma$ -rays with  $E_\gamma = 710 \pm 15, 1023 \pm 5, 1438 \pm 5, 2120 \pm 60, 2880 \pm 10$  and  $3560 \pm 50$  keV are observed ([1957HU79](#), [1957MC35](#)). Upper limits to the transitions  $1.74 \rightarrow \text{g.s.}$  and  $3.59 \rightarrow 1.74$  are  $2$  and  $3\%$  ([1964SI03](#)). See also ([1965SE1F](#)).



The summed proton energy spectrum, observed at  $E_p = 150$  to  $460$  MeV shows peaks corresponding to removal of an  $l \neq 0$  proton at  $Q = -6.7, -11.9$  and  $-17.1$  MeV; for removal of an  $l = 0$  proton,  $Q = -30.5$  MeV ([1966TY01](#)). See also ([1962DI1A](#), [1963BE42](#), [1964BA1C](#), [1964LI1D](#), [1965RI1A](#)) and  ${}^9\text{Be}$ .



Deuteron groups have been observed corresponding to eleven excited states of  ${}^{10}\text{B}$ : see Table [10.22](#). The absence of the groups (upper limit to intensity  $1.5\% - 4\%$ ) corresponding to the  $1.74$  and  $5.17$  MeV states is good evidence of their  $T = 1$  character ([1962AR02](#)). See also ([1959TO1A](#), [1962SL03](#)).



See ([1963HO19](#)).



See ([1956BO25](#), [1964ST1K](#)).



See ([1961AN07](#)).

Table 10.22:  $^{10}\text{B}$  levels from  $^{10}\text{B}(\text{p}, \text{p}')$ <sup>a</sup>,  $^{10}\text{B}(\text{d}, \text{d}')$  and  $^{12}\text{C}(\text{d}, \alpha)^{10}\text{B}$ 

$E_x$ <sup>b</sup> (MeV)	Relative Intensities <sup>b</sup>		$E_x$ <sup>c</sup> (p, p')	$E_x$ <sup>d</sup> (d, d')	$E_x$ <sup>d</sup> $^{12}\text{C}(\text{d}, \alpha)$	$E_x$ <sup>e</sup> (MeV)	$\Gamma$ <sup>e</sup> (keV)
	(p, p')	(d, d')					
0	100	100	0				
0.717 <sup>f</sup>	4	7	0.72	0.72	0.72		
1.739	1	< 0.15	1.74	g	g		
2.152	3	4	2.15	2.15	2.15		
3.583	7	7	3.58	3.58	3.58		
4.771			4.77	4.77	4.77		
			5.11	5.11	5.11		
			5.16	g	g		
				5.18		110 ± 10 <sup>h</sup>	
					5.92	< 5	
					6.03	< 5	
					6.13	< 5	
					6.55	25 ± 5	
					7.00	95 ± 10	
					7.48 <sup>i</sup>	90 ± 15	

<sup>a</sup> See also (1964JA03, 1965HA17).

<sup>b</sup> (1953BO70): energies ±5 keV. Relative intensities at  $\theta = 90^\circ$ ,  $E = 7.6$  MeV.

<sup>c</sup> (1962AR02): energies ±10 keV.  $E_p = 10.0$  MeV,  $\theta_L = 90^\circ, 120^\circ$ .

<sup>d</sup> (1962AR02): energies ±10 keV.  $E_d = 10.0$  MeV,  $\theta_L = 60^\circ, 90^\circ$ .

<sup>e</sup> (1964AR04):  $E_p$  and  $E_d = 11$  MeV,  $\theta_L = 45^\circ, 60^\circ$ ; energies ±10 keV.

<sup>f</sup> 719 ± 1.6 keV (1952CR30), 718 ± 5 keV (1954DA20).

<sup>g</sup> Absent: intensity of  $E_x = 5.17$  MeV group < 1% relative to average of  $T = 0$  levels.

<sup>h</sup> c.m. width (1962AR01).

<sup>i</sup> Relative intensity in (d, d') < 2%.



The half life is  $19.48 \pm 0.05$  sec ([1962EA02](#)),  $19.27 \pm 0.08$  sec ([1963BA52](#)):  $E_{\beta^+}(\text{max}) = 1.865 \pm 0.015$  MeV ([1963BA52](#)). The  $\beta^+$  decay is to the first two excited states of  $^{10}\text{B}$ : relative transition probabilities to the 0.72, 1.74 and 2.15 MeV levels are  $98.4/1.65 \pm 0.2/ < 0.1$  ([1953SH38](#)):  $ft$  (from  $\tau_{1/2} = 19.41$  sec,  $Q_m$  above) are  $1.0 \times 10^3$  and  $2.2 \times 10^3$  to the 0.72 and 1.74 MeV states, respectively ([1966BA1A](#)).



See ([1951SH63](#)).



At  $E_p = 19$  MeV, the ground state and the first four excited states have been observed. From the angular distributions, analyzed by PWBA,  $l_n = 1$ ,  $\theta^2 = 0.011, 0.029, 0.011$  and  $0.0031$  ([1961LE1A](#), [1963LE03](#)). At  $E_p = 155$  MeV, deuteron groups are reported to states at 0, 2.0, 5.0, 6.9 and 11.4 MeV ([1963BA2F](#)). See also ([1956RE04](#), [1960NE1C](#), [1961CL09](#), [1964SH07](#)).



See ([1960VL05](#), [1963OG1A](#)).



Reported levels are listed in Table 10.23. No evidence is found for previously reported levels at  $E_x = 2.86, 5.58$  and  $6.40$  MeV ([1963GA03](#), [1965GO05](#)). See also ([1965RO01](#)). The apparent absence of the  $E_x = 5.18$  MeV group is ascribed partly to its breadth and partly to its presumptive two-excited nucleon character ([1965GO05](#)). The angular distributions of alpha particles corresponding to the lower states indicate strong direct interactions,  $l = 1$  at  $E(^3\text{He}) = 2.3$  to  $5.5$  MeV ([1960TA12](#), [1965FO06](#): see also ([1965GO05](#))); the distributions of  $\alpha_2$  ( $E_x = 1.74$  MeV) vary strongly with energy ([1965FO06](#)). Alpha-gamma coincidence studies yield  $\Gamma_\gamma/\Gamma < 4 \times 10^{-4}$  for  $^{10}\text{B}^*(4.77)$ ,  $\approx 1$  for  $^{10}\text{B}^*(5.17)$  ([1965RO01](#)). See also ([1961BO33](#), [1964EL1B](#), [1964GA1B](#)). The  $\tau_m$  of the 1.74 MeV state is  $0.15 \pm 0.02$  psec ([1965LO04](#)).

Table 10.23:  $^{10}\text{B}$  levels from  $^{11}\text{B}({}^3\text{He}, \alpha)^{10}\text{B}$

$E_x$ <sup>a</sup> (MeV $\pm$ keV)	$d\sigma/d\Omega$ <sup>b</sup> (mb/sr)	$E_x$ <sup>c</sup> (MeV $\pm$ keV)	$d\sigma/d\Omega$ <sup>d</sup> (mb/sr)
0	0.69	0	0.30
$0.725 \pm 10$	0.38	$0.718 \pm 7$	0.09
$1.741 \pm 10$	0.29	$1.744 \pm 7$	0.06
$2.149 \pm 10$	0.34	$2.157 \pm 6$	0.13
$3.594 \pm 15$	0.26	$3.587 \pm 6$	0.12
		$4.777 \pm 5$	0.09
		$5.114 \pm 5$	0.10
		$5.166 \pm 5$	0.49
		$5.923 \pm 5$	0.09
		$6.028 \pm 5$	0.09
		$6.131 \pm 5$	0.09
		$6.573 \pm 8$	0.06
		$7.475 \pm 10$	0.30
		$7.567 \pm 10$	0.05

<sup>a</sup> (1960TA12):  $E({}^3\text{He}) = 5.2$  MeV.

<sup>b</sup> At  $\theta_{\max}$  (1960TA12).

<sup>c</sup> (1965GO05).

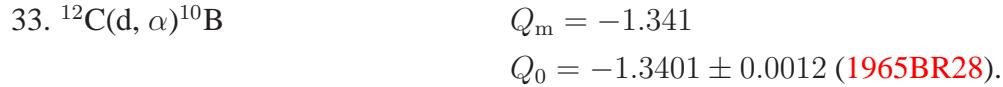
<sup>d</sup>  $E({}^3\text{He}) = 4.96$  MeV,  $\theta = 60^\circ$ : see (1965GO05) for other angles.



Not reported.



See (1961CL09, 1964BA1C).



Alpha groups have been observed to all of the known states of  $^{10}\text{B}$  up to 5.1 MeV: the intensity of the  $\alpha_2$ -group to the  $0^+$ ;  $T = 1$  state at 1.74 MeV is usually sharply reduced ([1957EL12](#), [1961PE09](#), [1961YA08](#), [1962AR02](#), [1963AL16](#), [1963YA1B](#), [1965BA06](#), [1965VO1B](#), [1966HA09](#)); but its yield varies appreciably with small energy changes ([1963AL16](#)). The identification of the  $\alpha$ -groups to the 5.1 MeV states indicates that the  $T = 1$  state at 5.17 MeV is not excited at  $E_d = 10$  MeV ([1962AR02](#)). See Table 10.22. Angular distributions of the  $\alpha_0$ -,  $\alpha_1$ -,  $\alpha_3$ - and  $\alpha_4$ -groups have been determined for  $E_d = 9.2$  to 19.7 MeV ([1961YA08](#), [1963YA1B](#), [1965BA06](#)). See also ([1958CA1F](#), [1959HE1C](#), [1961GA13](#), [1963JA03](#), [1963PE07](#), [1964BA54](#)).



See ([1962ZA01](#), [1964BR1L](#)).



At  $E_p = 12.2$  MeV,  $\alpha$ -groups are observed to the ground state and to the first three excited states ([1963PE07](#)). See also ([1962HA1F](#)).



For reaction (a) see ([1959AJ76](#)); for reaction (b) see ([1961CL09](#)).

Table 10.24: Energy levels of  $^{10}\text{C}$ 

$E_x$ (MeV $\pm$ keV)	$J^\pi$	$\Gamma$ (keV)	Decay	Reactions
g.s.	(0 $^+$ )	$\tau_{1/2} = 19.41 \pm 0.04$ sec	$\beta^+$	1, 2, 4, 5
3.360 $\pm$ 17			$\gamma$	2, 5
5.6 $\pm$ 100				2, 5
(7.2 $\pm$ 200)				5
10.2 $\pm$ 200		$\approx 1500$		5

$^{10}\text{C}$   
(Figs. 21 and 22)

GENERAL: See ([1960TA1C](#), [1962IN02](#), [1964GR1J](#), [1964VO1B](#)).

*Mass of  $^{10}\text{C}$ :* From the  $Q$ -value of the  $^{10}\text{B(p, n)}^{10}\text{C}$  reaction ([1961TA12](#):  $Q_0 = -4.393 \pm 0.025$  MeV) and the  $\beta^+$  end-point energy ([1963BA52](#):  $E_{\beta^+}(\text{max}) = 1.865 \pm 0.015$  MeV), the mass excess of  $^{10}\text{C}$  is  $15.658 \pm 0.013$  MeV, based on  $^{12}\text{C} \equiv 0$  ([1965MA54](#)).



The decay is complex. See  $^{10}\text{B}$ .



The ground state threshold occurs at  $E_p = 4.835 \pm 0.025$  MeV ([1961TA12](#)). A sharp rise in the cross section curve at  $E_p = 8.55$  MeV is attributed to a state in  $^{10}\text{C}$  at  $E_x = 3.38 \pm 0.03$  MeV ([1963EA01](#)). This state has also been observed directly as a neutron group ( $E_x = 3.34 \pm 0.2$  MeV) at  $E_p = 17.2$  MeV. Wide or unresolved states at  $E_x \approx 5.1$  MeV may also be indicated ([1954AJ32](#)). The  $\gamma$ -decay of the first excited state has been observed:  $E_x = 3.35 \pm 0.02$  MeV ([1965SE1F](#)). See also ([1965VA23](#)) and  $^{11}\text{C}$ .



See ([1964OS1A](#)).



See ([1959OC07](#)).



At  $E_p = 155$  MeV, triton groups are observed to the ground state of  $^{10}\text{C}$  and to excited states at  $3.4 \pm 0.1$ ,  $5.6 \pm 0.1$ ,  $(7.2 \pm 0.2)$ , and  $10.2 \pm 0.2$  MeV. The 10.2 MeV state has a width of  $\approx 1.5$  MeV ([1965BA17](#)). See also ([1960RA12](#), [1961RA1B](#), [1962CO17](#)).

## References

(Closed July 01, 1965)

References are arranged and designated by the year of publication followed by the first two letters of the first-mentioned author's name and then by two additional characters. Most of the references appear in the National Nuclear Data Center files (Nuclear Science References Database) and have NNDC key numbers. Otherwise, TUNL key numbers were assigned with the last two characters of the form 1A, 1B, etc. In response to many requests for more informative citations, we have, when possible, included up to ten authors per paper and added the authors' initials.

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